Particle Physics II (LPHY2133)

Andrea Giammanco, UCL



Putting things together

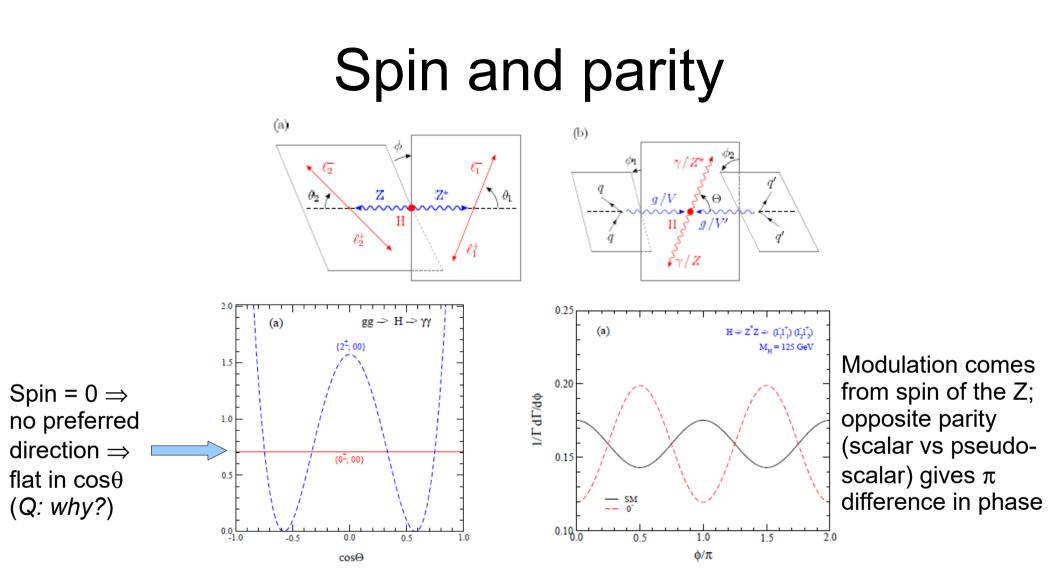
Andrea Giammanco

Complementary information

- The point about studying several final states is not only about getting more significance in the combination
 - Although this was an important consideration until 2012
- Different final states give complementary information
- Once you discover a new particle you also want to know which particle you have discovered
 - Those analyses were designed for the Higgs boson, but a different new particle may have passed the same selection
- Next slides are about how our understanding of this new particle has formed, based on the available data
- (For another historical case, you can compare to how the J/ψ 's identity was understood in 1974: link)

Spin of the new particle

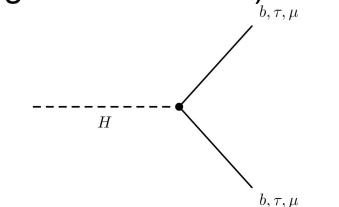
- From the fact that it decays into 2 γ, 2 Z, 2 W (all spin=1 particles), we know:
 - Spin cannot be fractional (\Rightarrow it's a boson)
 - Decay into two spin-1 particles limits the spin to 0, 1, 2
 - (spin >2 only if large orbital angular momentum of decay products; not impossible but strongly disfavoured)
- The decay into two real γ 's also excludes spin = 1
 - Because the photon is massless, hence it has only two polarization states (Landau-Yang theorem, ref.: link)
- Angular distributions of the decay products are affected by the spin and parity of the intermediate resonance

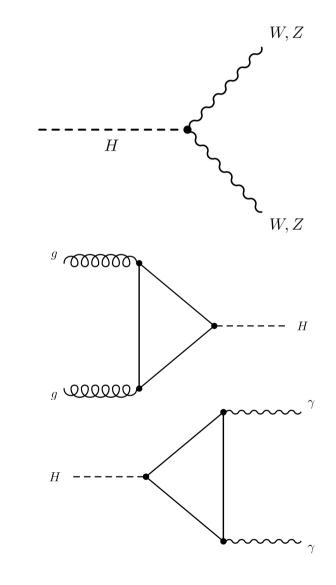


- Figures taken from a nice ref.: link; another good theory source: link
- Scalar hypothesis (J^P=0⁺) confirmed by ATLAS and CMS in Run-1, combining several angles in MVA discriminants in ZZ, γγ, WW
- Most precise results from $H \rightarrow ZZ \rightarrow 4I$ (*Q: say at least two reasons*)

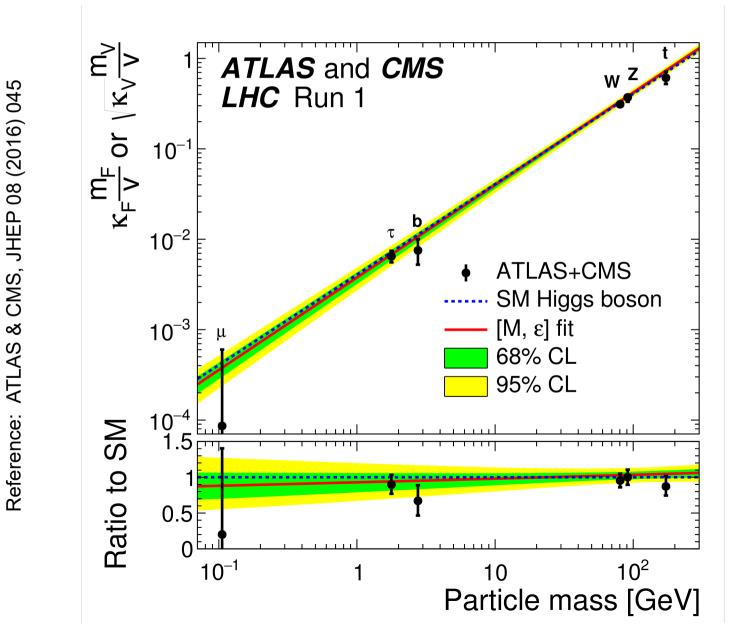
Couplings

- It decays to WW and ZZ ⇒ it couples to W and Z, as expected by the Higgs mechanism to explain their mass
- It is produced by gg→H and decays to γγ ⇒ *indirect* proof that it couples also to fermions (or at least the top quark)
- It decays to ττ and bb ⇒ direct proof that it couples also to fermions (or at least 3rd generation ones)





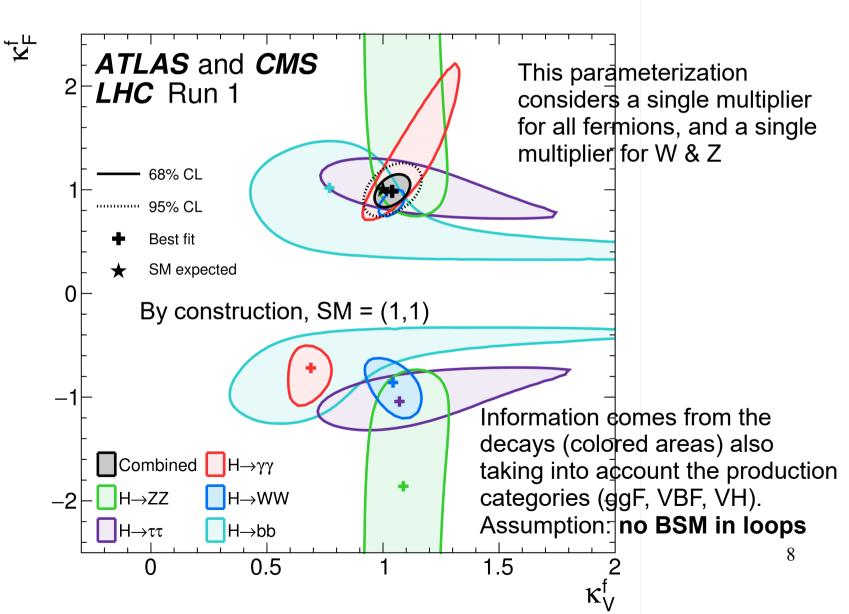
Couplings: quantitative tests of the SM's Higgs mechanism



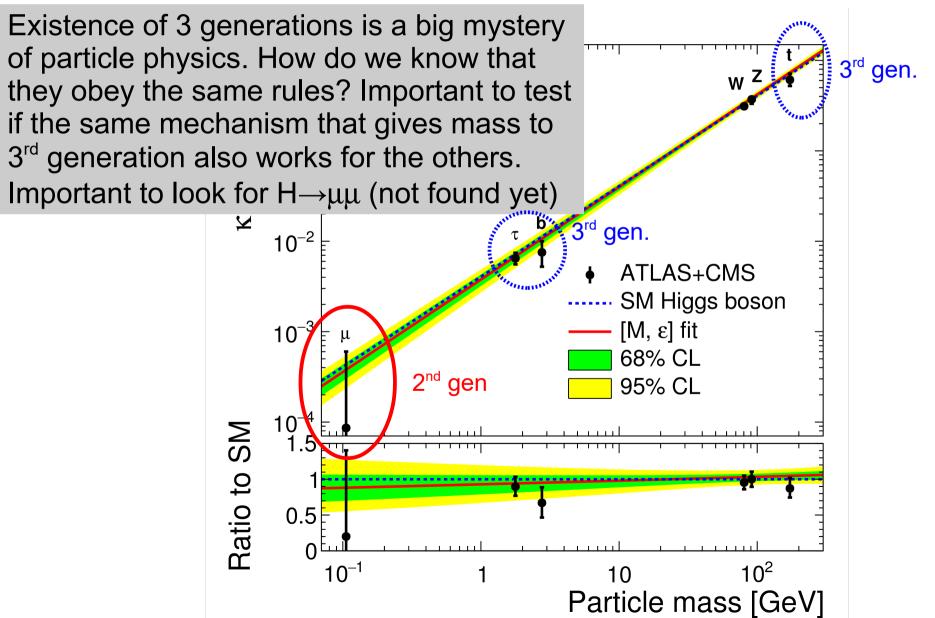
Fermion (k_F) and boson (k_V) coupling multipliers

Quiz:

- Explain ZZ
- Explain $\tau\tau$
- Explain $\gamma\gamma$

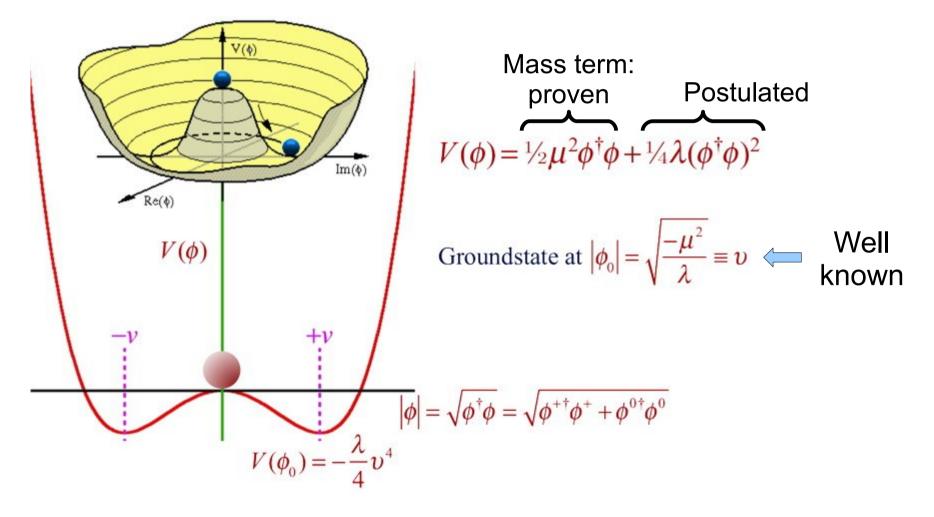


What's next: are all generations the same?



9

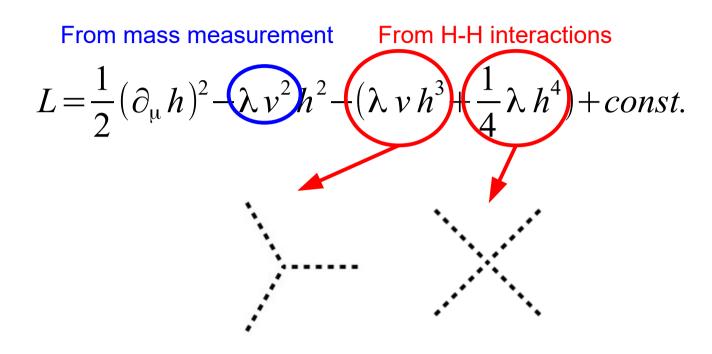
What's next: testing the V(ϕ) slope



Picture from https://web2.ph.utexas.edu/~coker2/index.files/gaugef.htm

10

The Higgs self-interactions



Parameter λ is *indirectly* known from our knowledge of v and m_u.

Reactions of the kind HH \rightarrow H, H \rightarrow HH and HH \rightarrow HH would allow a model-independent knowledge of the Higgs potential, to be compared to the shape dictated by the SM.

This is one of the research directions at CP3.



The top-Higgs connection

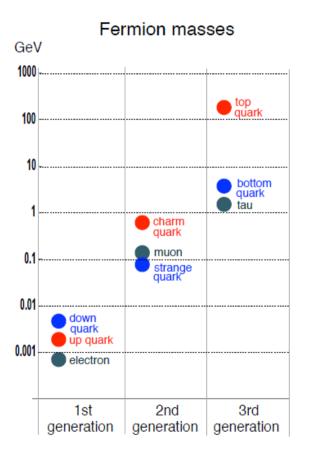
Andrea Giammanco

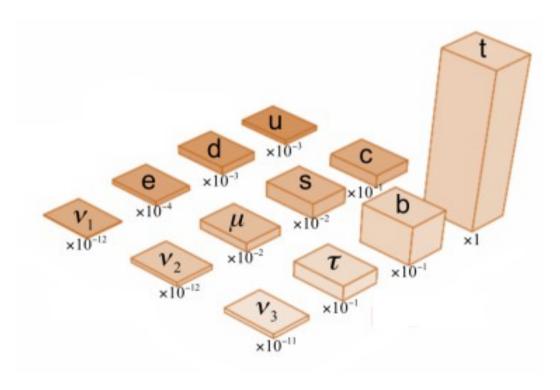
The flavour problem

FERMIONS matter constituents spin = 1/2, 3/2, 5/2,						
	Leptons spin =1/2			Quarks spin =1/2		
	Flavor	Mass GeV/c ²	Electric charge	Flavor	Approx. Mass GeV/c ²	Electric charge
	\mathcal{V}_{L} lightest neutrino* e electron	(0−2)×10 ^{−9} 0.000511	0 -1	u _{up} d down	0.002 0.005	2/3 1/3
	$\mathcal{V}_{\mathbf{M}}$ middle neutrino* μ muon	(0.009–2)×10 ^{–9} 0.106	0 -1	C charm S strange	1.3 0.1	2/3 -1/3
	\mathcal{V}_{H} heaviest neutrino* au _{tau}	(0.05–2)×10 ^{–9} 1.777	0 -1	t _{top} b _{bottom}	173 4.2	2/3 -1/3

Nicely arranged in three generations. But we don't know why. Image source: CPEP

No clear pattern





From Gilad Perez

Where is the top mass coming from?

Z=- Fre Friv titte +h.c. + (Y; Y; Y; Ø+ h.c. + $D_{\mu}\phi l^2 - V(\phi)$

15

Where is the top mass coming from? $\phi = v + h \qquad L \supset \frac{y_t}{\sqrt{2}} (v \bar{\psi}_t \psi_t + \bar{\psi}_t \psi_t h)$

In the SM, all fermion mass terms come from the Electro-Weak Symmetry Breaking.

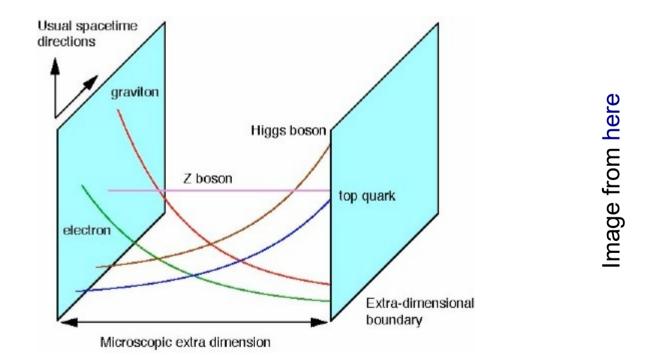
The Vacuum Expectation Value, v, is well known from other fundamental parameters of the SM.

We have made some progress: we now know that the fermion mass hierarchy is a mere reflection of the hierarchy in Yukawa coupling strengths. But no explanation for that hierarchy of Yukawa couplings

16

An example of a deeper explanation

• Randall-Sundrum mechanism (string-inspired):



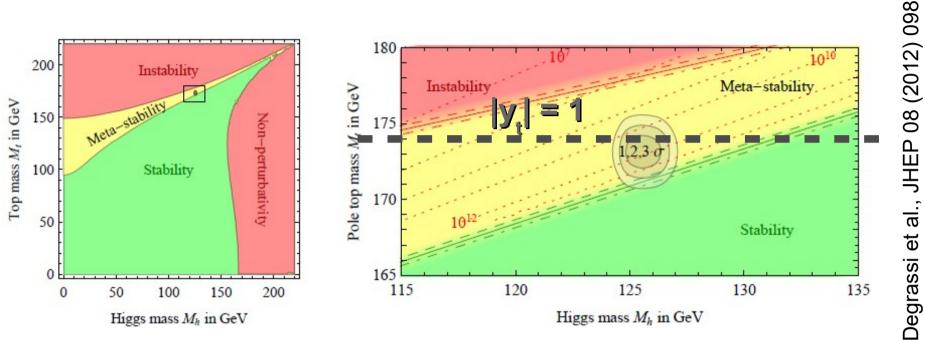
Coupling strengths come from the wavefunction shapes and their overlaps in the warped dimension. Other models assume a special role of the top quark.

Yukawa coupling of the top

$$L \supset \frac{y_t v}{\sqrt{2}} \overline{\psi}_t \psi_t \equiv M_t \overline{\psi}_t \psi_t$$

- Fill actual numbers in:
 - $v = (\sqrt{2}G_F)^{-1/2} = 246.2196(1) \text{ GeV} (G_F = 1.166 378 7(6)x10^{-5} \text{ GeV}^{-2})$
 - $M_{t}^{CMS} = 172.44 \pm 0.49 \text{ GeV}$ (CMS coll., Phys. Rev. D 93 (2016) 072004)
 - $M_{t}^{\text{Tevatron}} = 174.30 \pm 0.65 \text{ GeV}$ (CDF&D0 coll., FERMILAB-CONF-16-298-E)
- We get $y_t^{CMS} = 0.990 \pm 0.003$, $y_t^{Tev} = 1.001 \pm 0.004$
 - Closeness to 1 is interesting, for an adimensional parameter
 - Pure chance, or does it reflect something deep?
 - The SM offers no explanation (apart from pure chance)

Stability of the Universe



To understand the red areas, see also Sec.2.5 of "The Higgs Hunter's Guide"

- This study assumes SM validity up to the Planck scale; and, in the SM, m₁ and m₁ are free parameters
- Under these assumptions, conspiracy of top and Higgs makes our Universe sit on the thin line between stability and instability
- We don't know if there is a deep reason for that

prolific literature: link)

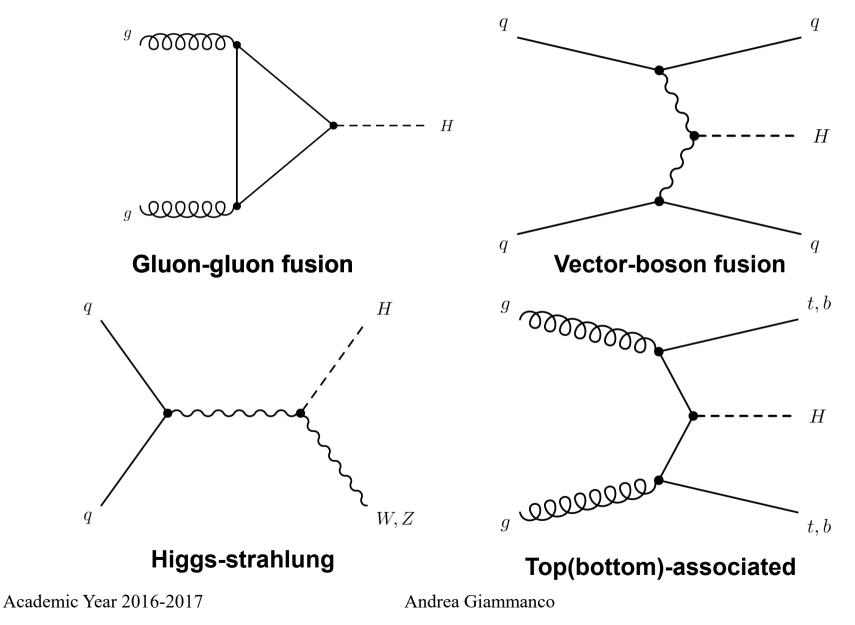
σ

stimulated

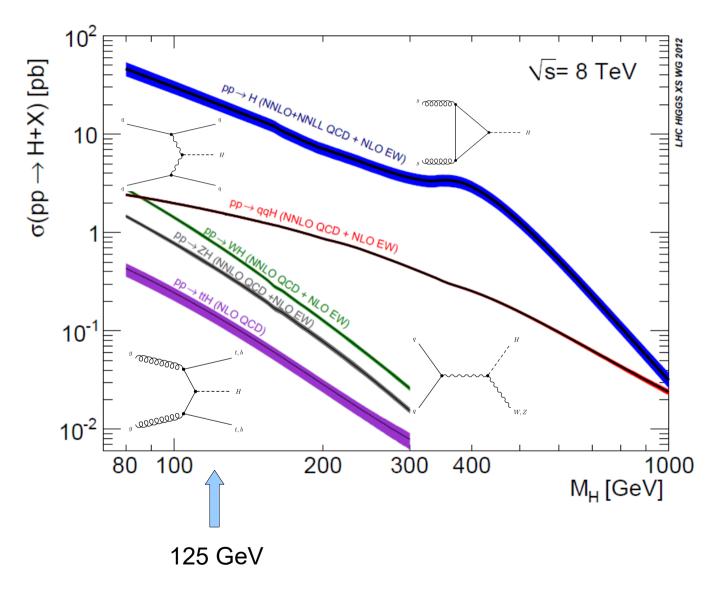
 $y_{t} = \sqrt{2} M_{t} / v ?$

- Crucial test of the SM
 - If the equality is not exact, we prove that fermion masses are not (only) due to the Higgs mechanism
- To answer this question, we need to:
 - Measure the mass of the top quark precisely
 - (Not in the scope of this lecture; see my slides here)
 - Measure the Yukawa coupling through observables independent from the top mass
 - The most direct is the ttH cross section
 - Note: ttH has not even been observed (*) yet

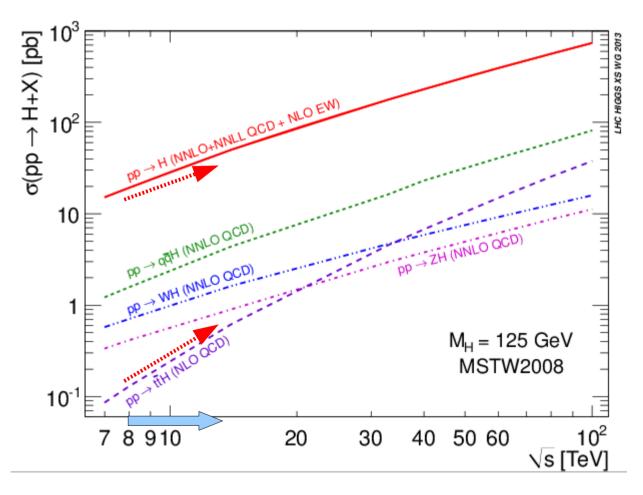
How to produce a Higgs boson



Cross sections



As function of collision energy

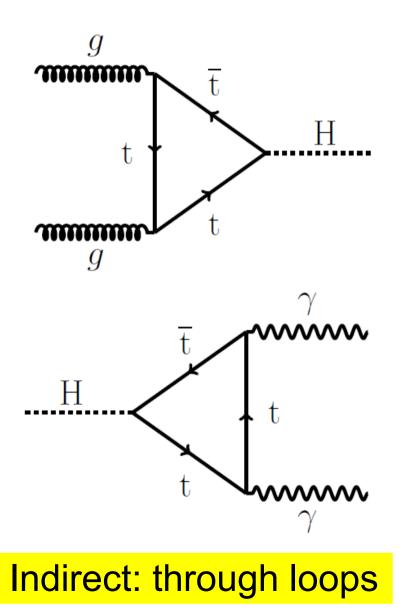


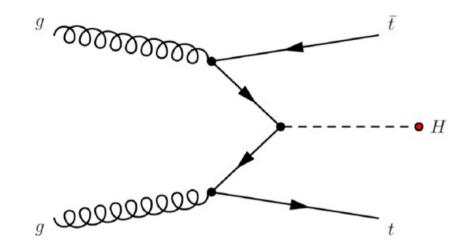
The larger the energy, the easier is ttH From 8 to 13 TeV, 4x increase in cross section

Academic Year 2016-2017

Andrea Giammanco

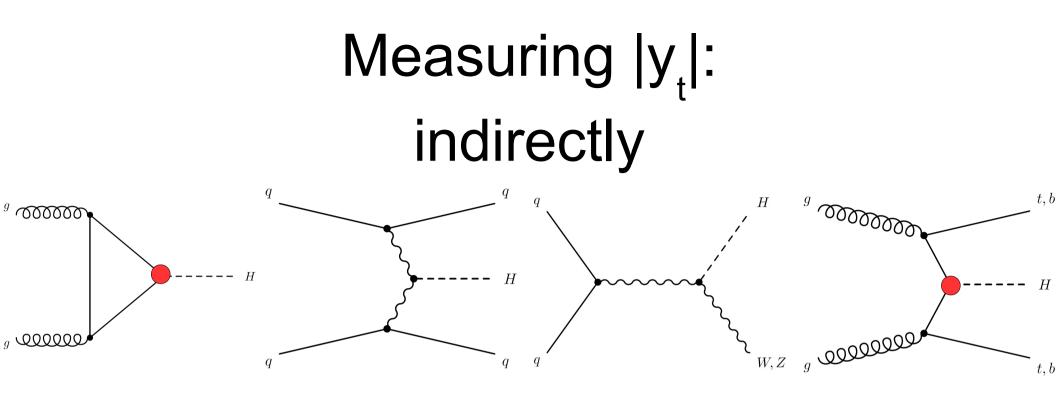
Measuring |y_t|





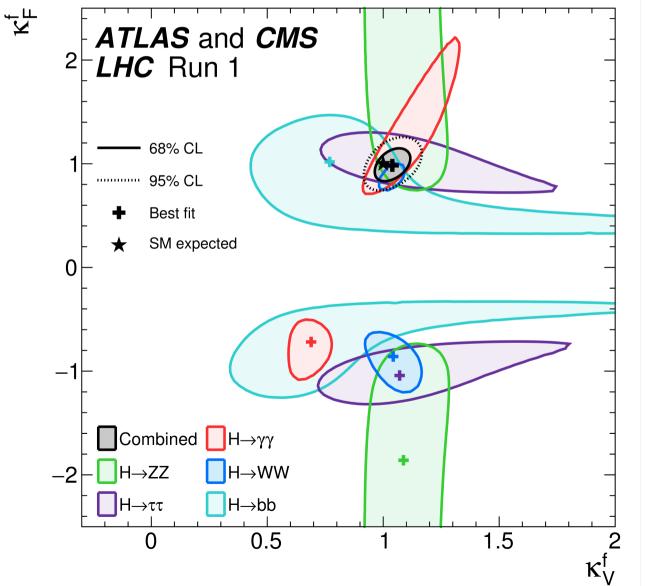
Direct: at tree level

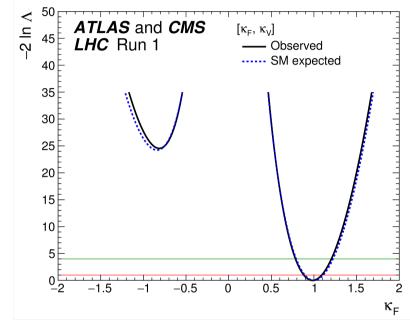
General consideration: when sensitivity is induced by loops, one needs to rely more on some model assumptions (e.g., what particles run in the loop)



- Joint ATLAS+CMS run-1 Higgs properties paper (*) combines many final states to extract constraints on several couplings, with several alternative parameterizations and assumptions
- Diagrams with indirect (loop) and direct (tree) sensitivity to the top-Higgs coupling are both considered but, at the current state, precision on this parameter is driven by the loops
 - (For sake of clarity, I will not elaborate on the role of ttH and tH in this global combination; explicit ttH search presented later)

Fermion (k_F) and boson (k_V) coupling multipliers

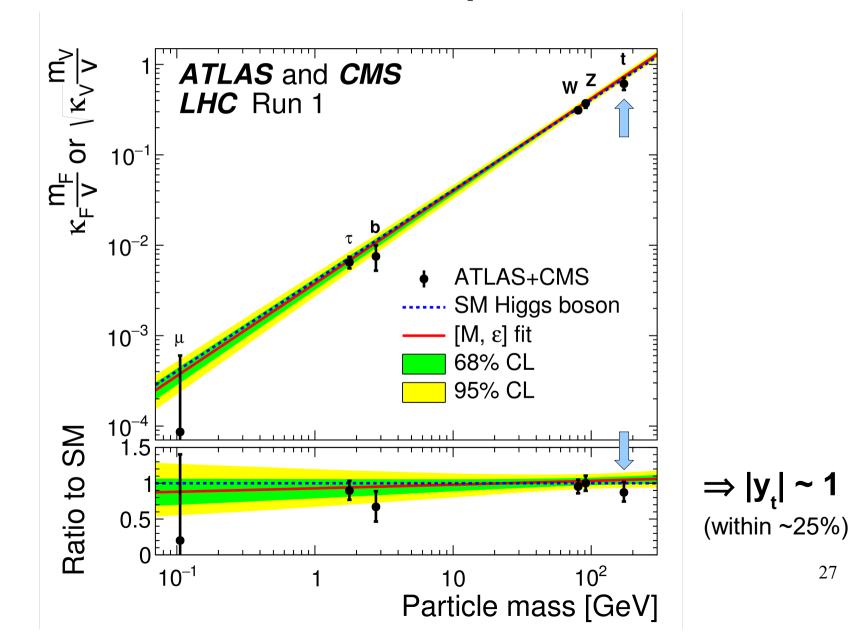




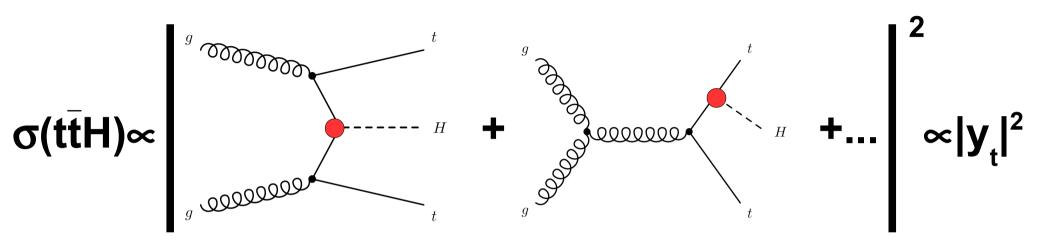
This parameterization considers a single multiplier for all fermions. $k_F \sim 1 \Rightarrow y_t \sim 1$ (within ~25%)

Assumption: no BSM in loops

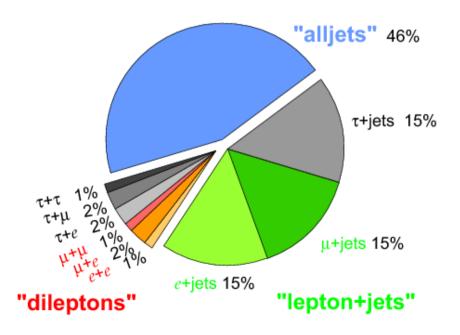
A test of the coupling-mass relationships



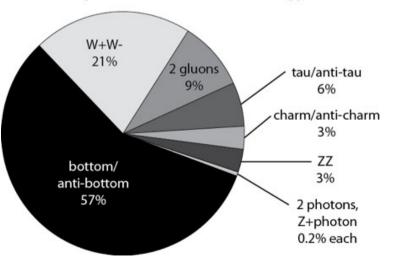
Measuring |y_t|: directly



Top Pair Branching Fractions



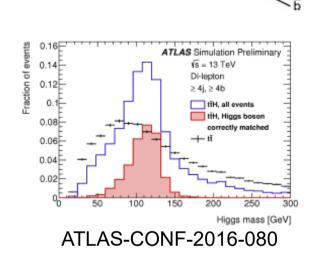
Decays of a 125 GeV Standard-Model Higgs boson



Searches for ttH: bb channel



- Largest BR (~60%)
- Large multiplicity of jets and b-tags
- Cons:
 - Overwhelming tt+jets background

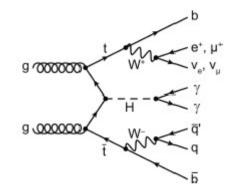


g 000000

,00000

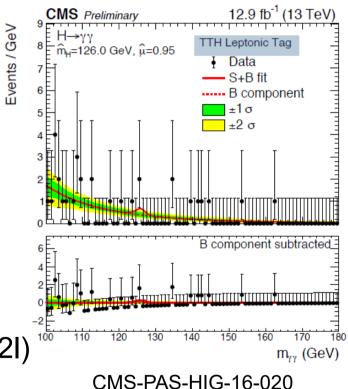
- Heavy flavour component of bkg (ttbb, ttcc) is poorly constrained
- Very large combinatorics of jet-parton associations
- Approaches:
 - Use several combinations of lepton / jet / b-tag multiplicities in simultaneous fit; it helps a lot in constraining bkg fractions
 - (MVA for jet-parton association, followed by) MVA for classification

Searches for tτH: γγ channel

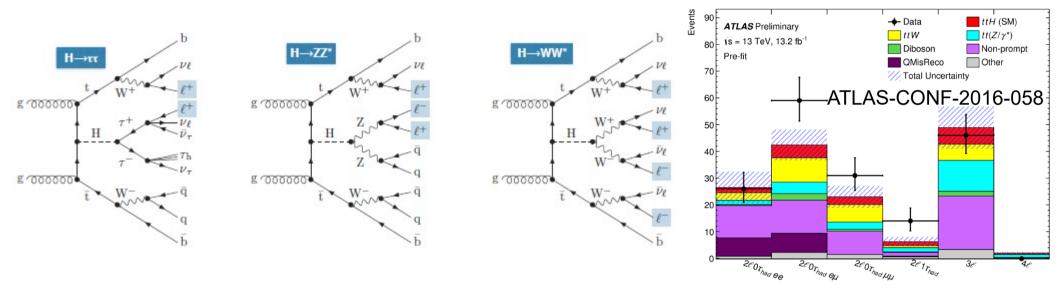


• Pros:

- High-resolution mass peak
- Background: smoothly falling mass spectrum
- Cons:
 - Small branching ratio (~0.2%)
- Approach:
 - Similar to standard γγ analysis
 - All possible tt final states are considered (0I,1I,2I)
 - Request two b-tagged jets

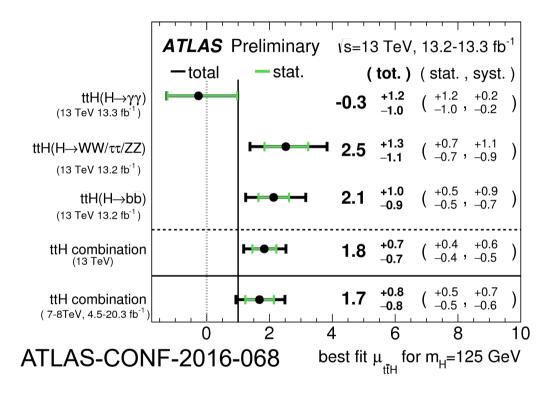


Searches for ttH: multi-lepton



- Very clean selections: 2 same-sign leptons, or ≥3 leptons
 - Target final states with at least one leptonic top decay and more leptons from H→ττ (6.3%), H→ZZ (2.6%), H→WW (21.5%)
- - Control regions are used to estimate fake rate from data
 - To estimate charge confusion: use $Z \rightarrow I^+I^+$ and $Z \rightarrow I^-I^-$ (*Q: how?*)

Searches for ttH: latest results (to be updated next Monday!)



CMS (13-15/fb):

• H→γγ: μ = 1.9^{+1.5} -1.2 CMS-PAS-HIG-16-020

• Multilepton: $\mu = 2.0^{+0.8}$ CMS-PAS-HIG-16-022

• H→bb: μ = -0.2±0.8 CMS-PAS-HIG-16-038

LHC Run-1 combination: • $\mu = 2.3^{+0.7}_{-0.6}$

~40% uncertainty on signal strength (μ) \Rightarrow ~20% on y_t ($\mu \equiv \sigma_{_{obs}}/\sigma_{_{exp}} \propto y_t^2 \Rightarrow \Delta \mu/\mu = 2\Delta y_t/y_t$)

Exam

- It will be on 23/06 (not 20/06) starting at 14:00
- Format:
 - Written "review" report of <10 pages in pdf format including pictures, references, etc., by 16/06
 - Oral exam (not a presentation!) on 23/06
 - Bonus if you do well at the mid-term evaluation next week
 - Bonus (up to 2 points in total) for the problem-solving evaluations during K.P.'s lectures

The written report

- Topics:
 - Muon collider: physics motivations
 - Muon collider: experimental challenges
 - New acceleration techniques
 - Dark matter searches (choose one type)
 - Free subject (but very well motivated)
- Purpose:
 - A short overview of the subject (in English or French)
- Structure:
 - Introduction: motivation + stating a "problem" + wider context
 - Present status of their studies and their relevance
 - Outlook: next steps and longer-term perspectives

